



Concurrent multiple impact in rigid bodies: Formulation and simulation

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Enoc 2005, Eindhoven, Netherlands, 7-12 august 2005

Outline

→ 1 – Motivations

1.1 – What is a concurrent multiple impact?

1.2 – Industrial motivations

1.3 – Main objectives

1.4 – Previous results

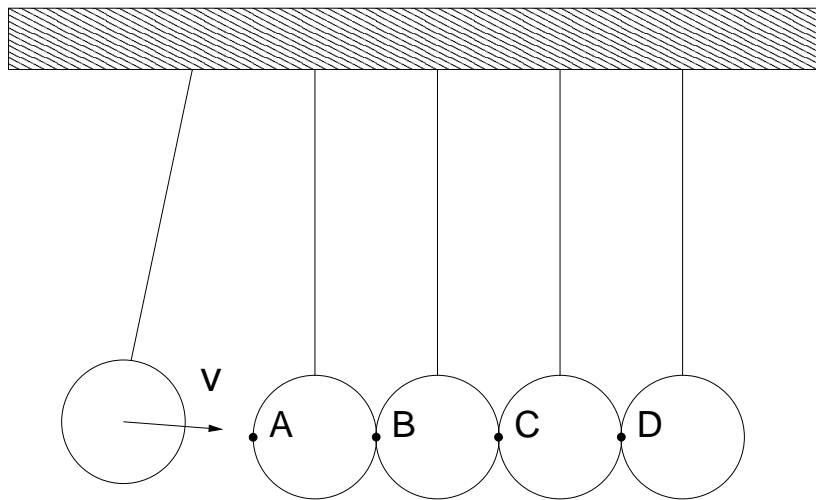
1.5 – Proposed approach

- 2 – Chain of balls
- 3 – Lagrangian systems
- 4 – Conclusion and perspectives

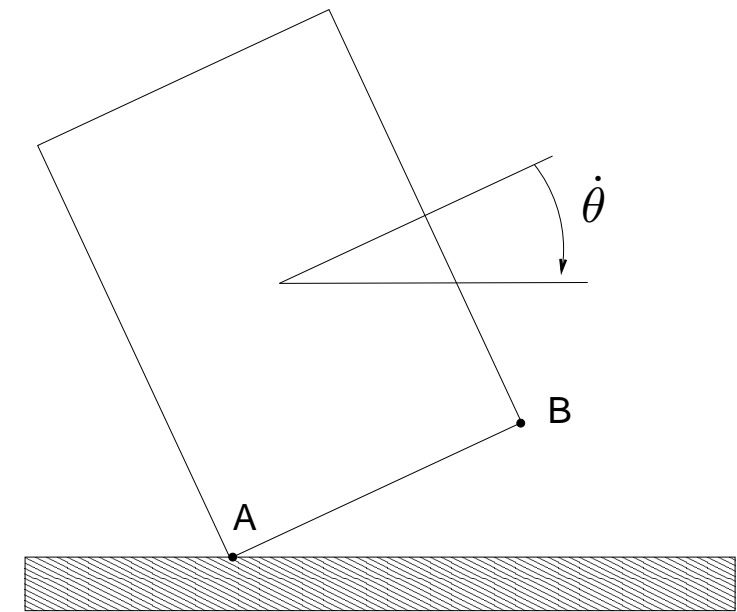
What is a concurrent multiple impact?

- ✱ A multiple impact in a multibody system may be defined as :
The occurrence of several impacts at the same instant on a rigid body.
- ✱ Academical examples :

Newton's cradle

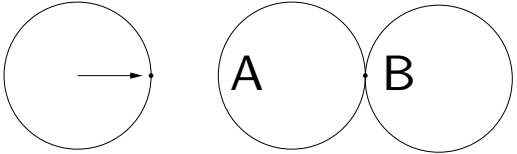
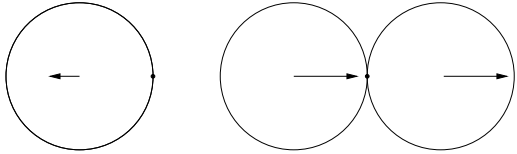
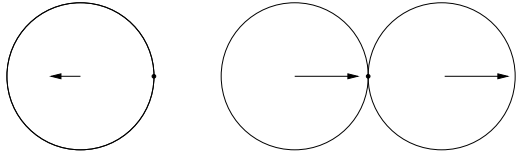
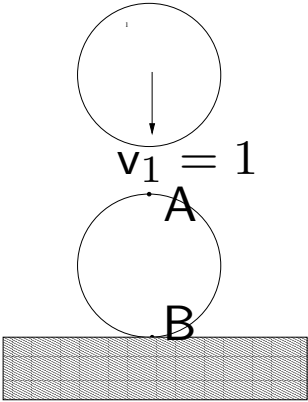
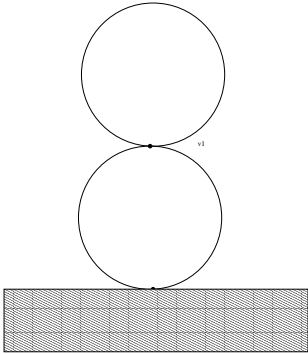
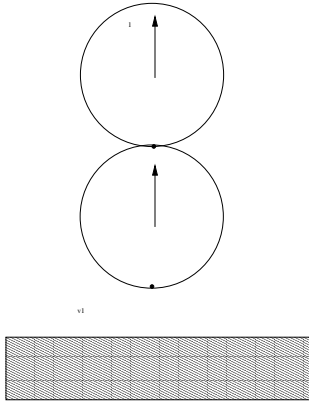


The rocking block



- ✱ Major difficulties:
 - The standard laws (Newton, Poisson, ...) do no longer apply correctly
 - The continuity with the respect to initial conditions is usually lost.

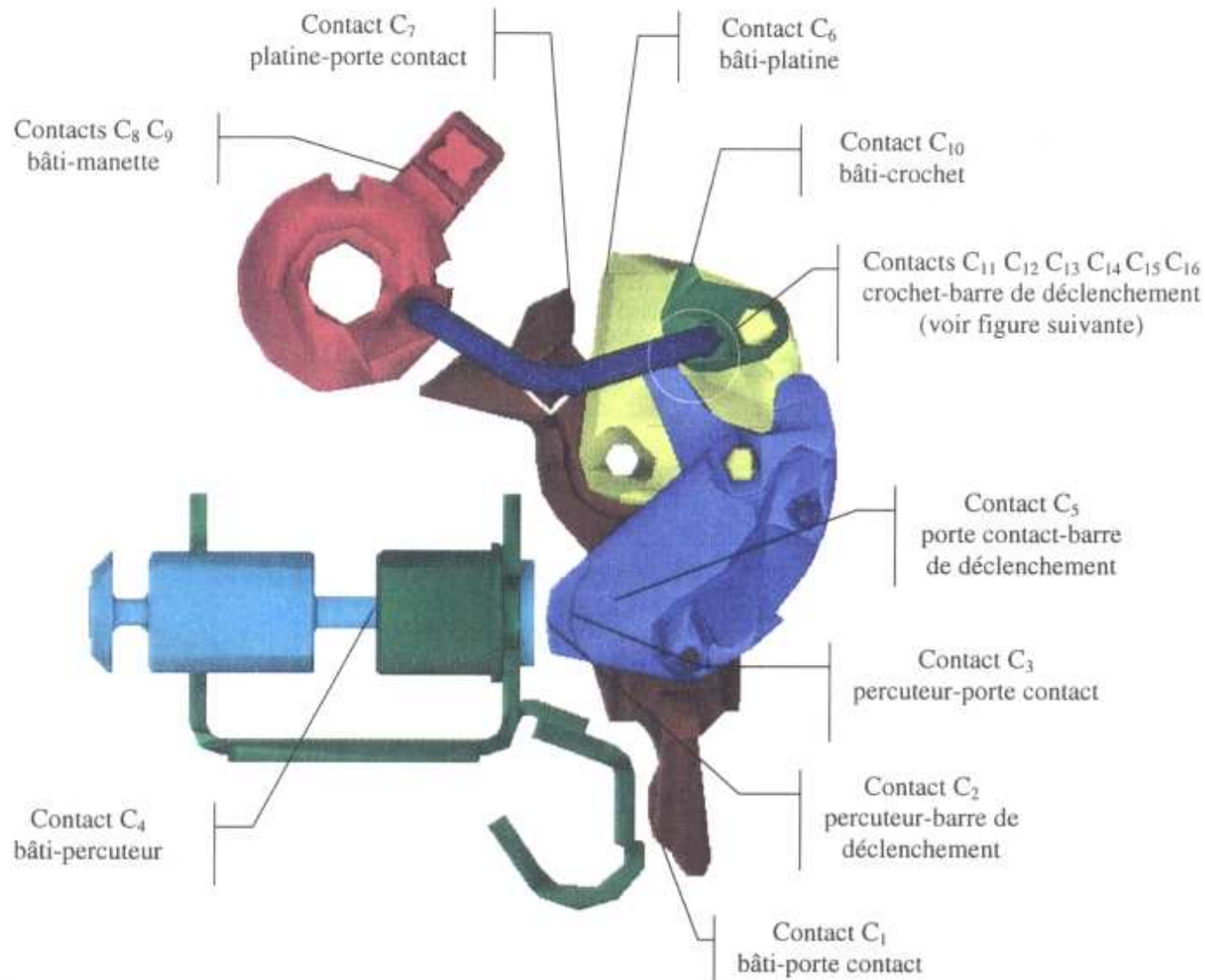
What is a concurrent multiple impact? Continued ...

Initial conditions	Newton	Poisson
 <p>$v_1 = 1$</p>	<p>$e_A = e_B = 1$</p>  <p>$v_1^+ = -1/3$ $v_2^+ = v_3^+ = 2/3$</p>	<p>$e_A = e_B = 1$</p>  <p>$v_1^+ = -1/3$ $v_2^+ = v_3^+ = 2/3$</p>
 <p>$v_1 = 1$</p>	<p>$e_A = 0, e_B = 1$</p>  <p>$v_1^+ = v_2^+ = 0$</p>	<p>$e_A = 0, e_B = 1$</p>  <p>$v_1^+ = v_2^+ \neq 0$</p>

Industrial motivations

Long term collaboration with Schneider Electric (M. Abadie)

Study of the stability of circuits breakers w.r.t. external impact excitations



Main objectives

Find a multiple impact law, i.e.

$$\dot{q}^+ = \mathcal{F}(\dot{q}^-, q, t)$$

with the following properties :

- ✱ single-valued mapping
- ✱ Ensure the fundamental principles of the mechanics of multibody system
 - Equations of motion
 - Unilateral constraints
 - Energetic balance
- ✱ Fit qualitatively and quantitatively the experiments
 - Measurable parameters with a physical meaning
- ✱ Efficient numerical treatment
 - Formulation in terms of a mathematical programm with constraints

Previous works

- ✱ **Sequential approach [Han & Gilmore , 1993]**
 - Heuristic choice of solutions
 - Existence and uniqueness problems (closed-loop systems)
- ✱ **Generalization to multi-constrained systems**
(Newton (Moreau, 1988), Poisson (Glocker & Pfeiffer, 1995)).
 - Efficient numerical algorithm
 - Do not fit the experiments.
- ✱ **Introduction of an impulse ratio [Hurmuzlu, 2001]**
 - Lack of efficient numerical algorithm
 - Do not ensure the fundamental principles of mechanics
- ✱ **Non local formulation of multi-body systems with impacts [Frémond, 1995].**
 - Ensure the fundamental principles of mechanics
 - Lack of physical interpretation of the parameter
- ✱ **Geometric framework for impacting systems [Aeberhard & Glocker]**

Proposed approach

✱ Fundamental problem:

- Taking into account the impact propagation by wave effects.
- Incorporating a continuum mechanics model for the deformation behaviour

✱ Quasi-rigid body model: general assumption of the Hertz's contact

- Local deformation in the area of contact
- Quasi-static behavior laws (no inertia effects in the contact region)
- Negligible wave effects inside the bodies (focusing on bodies without special aspect ratios)

→ Multibody systems are modeled by a discrete dynamical systems of rigid bodies and unilateral contact force model (non-linear stiffness and damping)

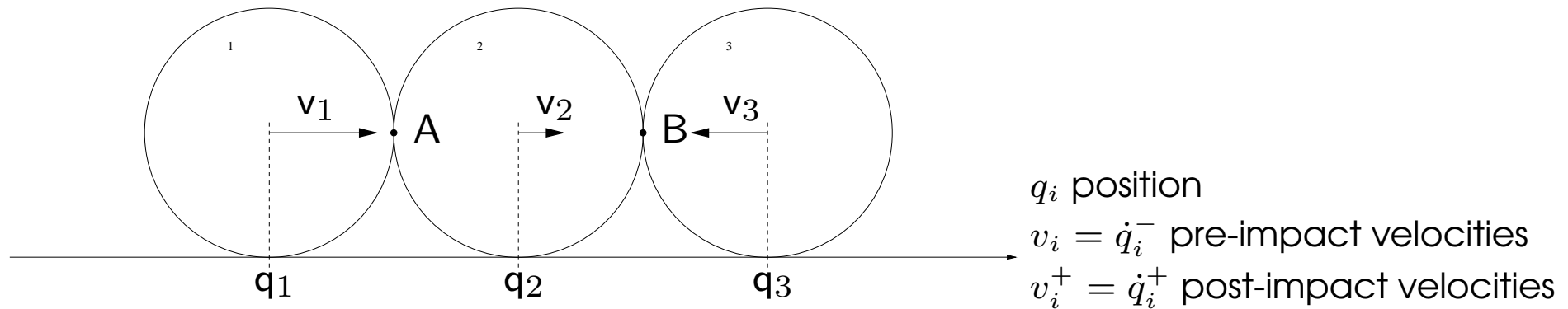
✱ A typical example of such systems: a chain of balls [Falcon & Laroche, 1998]

- Balls are modeled by discrete masses
- Contact interaction: Non linear Hertzian model

Outline

- ✓ 1 – Motivations
- 2 – Chain of balls
 - 2.1 – The unidimensional chain of 3 balls
 - 2.2 – The unidimensional chain of n balls
 - 2.3 – The pool: A bidimensional chain of balls
- 3 – Lagrangian systems
- 4 – Conclusion and perspectives

- The unidimensional chain of 3 balls



Equation of motion at the instant of impact ($v_2 = 0$) :

$$\begin{cases} m(v_1^+ - v_1) = -p_1 \\ m(v_2^+) = p_1 - p_2 \\ m(v_3^+ - v_3) = p_2 \end{cases} \quad (1)$$

Unilateral constraints

$$\begin{cases} v_1^+ \leq v_2^+ \leq v_3^+ \\ p_1 \geq 0, \quad p_2 \geq 0 \end{cases} \quad (2)$$

Choice of a parametrization

- ✱ One global energetic coefficient e :

$$(v_1^+)^2 + (v_2^+)^2 + (v_3^+)^2 = e(v_1^2 + v_3^2) \quad (3)$$

- ✱ One impulse ratio

$$\alpha = \frac{p_2}{p_1} \quad (4)$$

→ Parametrization of the solutions :

$$\begin{aligned} p_1 &= \frac{(1 + \sqrt{\Delta})m}{2(1 - \alpha + \alpha^2)} \\ p_2 &= \frac{(1 + \sqrt{\Delta})m\alpha}{2(1 - \alpha + \alpha^2)} \end{aligned} \quad (5)$$

$$\begin{aligned} \dot{q}_1^+ &= -\frac{1 + \sqrt{\Delta}}{2(1 - \alpha + \alpha^2)} + \dot{q}_1^- \\ \dot{q}_2^+ &= -\frac{(1 + \sqrt{\Delta})(\alpha - 1)}{2(1 - \alpha + \alpha^2)} \\ \dot{q}_3^+ &= -\frac{(1 + \sqrt{\Delta})\alpha}{2(1 - \alpha + \alpha^2)} \end{aligned} \quad (6)$$

with $\Delta = -1 + 2e - 2\alpha e + 2\alpha + 2\alpha^2 e - 2\alpha^2$.

Evaluation of e and α

- ✱ Bounds imposed by $\Delta > 0$ and unilateral constraints

$$\frac{1}{2} \leq \alpha \leq \frac{2e - 2 - \sqrt{6e - 2}}{e - 3}, \quad \frac{1}{3} \leq e \leq 1 \quad (7)$$

- ✱ Evaluation by a regularized system with Hertzian springs.

$$\begin{cases} m\ddot{\delta}_1 = -2f_1(\delta_1) + f_2(\delta_2) \\ m\ddot{\delta}_2 = -2f_2(\delta_2) + f_1(\delta_1) \\ 0 \leq \mathbf{f} \perp \mathbf{f} - \mathbf{K}(\boldsymbol{\delta}) \cdot \boldsymbol{\delta} \geq 0 \end{cases} \quad (8)$$

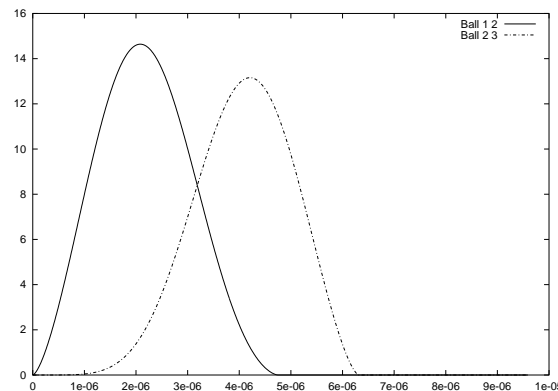
$$\delta_i = q_{i+1} - q_i \text{ Indentation, } K(\boldsymbol{\delta}) = \begin{bmatrix} k_1 \delta_1^{1/2} & 0 \\ 0 & k_2 \delta_2^{1/2} \end{bmatrix} \quad \kappa = \frac{k_1}{k_2} \quad (9)$$

- ✱ Evaluation of the impulse ratio

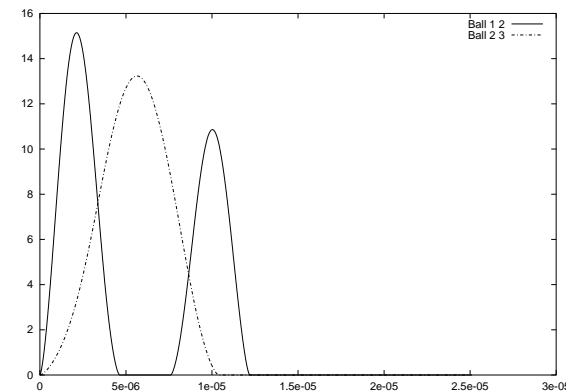
$$\alpha = \frac{\int_0^{t_f} f_1(t) dt}{\int_0^{t_f} f_2(t) dt} \quad (10)$$

- Evaluation of e and α (Continued ...)

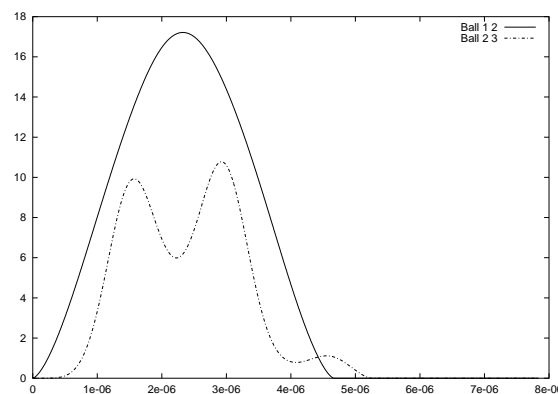
Numerical integration of 3 balls chain. Hertzian spring contact.



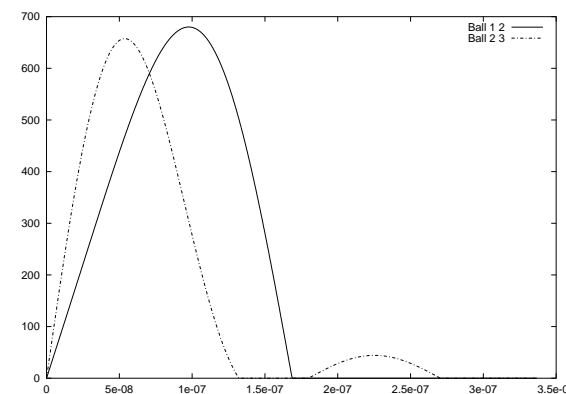
$$\kappa = 1, v_1 = 1, v_3 = 0$$



$$\kappa = 0.1, v_1 = 1, v_3 = -1$$



$$\kappa = 100, v_1 = 1, v_3 = 0$$



$$\kappa = 2.24, v_1 = 1, v_3 = -1$$

Forces between balls versus time.

→ No sequential or simultaneous process

The unidimensional chain of n balls

✱ Rigid body model of a unidimensional chain of n balls :

- Unknowns: n velocities \dot{q}_i , $n - 1$ impulses p_i
- Equations:
 - n equations of motion
 - 1 Energetic balance e
 - $n - 2$ impulse ratios, $\alpha_{ij} = p_i/p_j$

→ Unique determination of post-impact velocities and impulses

✱ Regularised quasi-rigid model:

- Visco-elastic contact model

$$f = K\delta^\nu + C\delta^{\nu-1}\dot{\delta} \quad (11)$$

- Evaluation of the impulse ratios

$$\alpha_{\gamma\beta} = \frac{\int_0^{t_f} f_\gamma(t) dt}{\int_0^{t_f} f_\beta(t) dt} \quad (12)$$

The unidimensional chain of n balls (Continued ...)

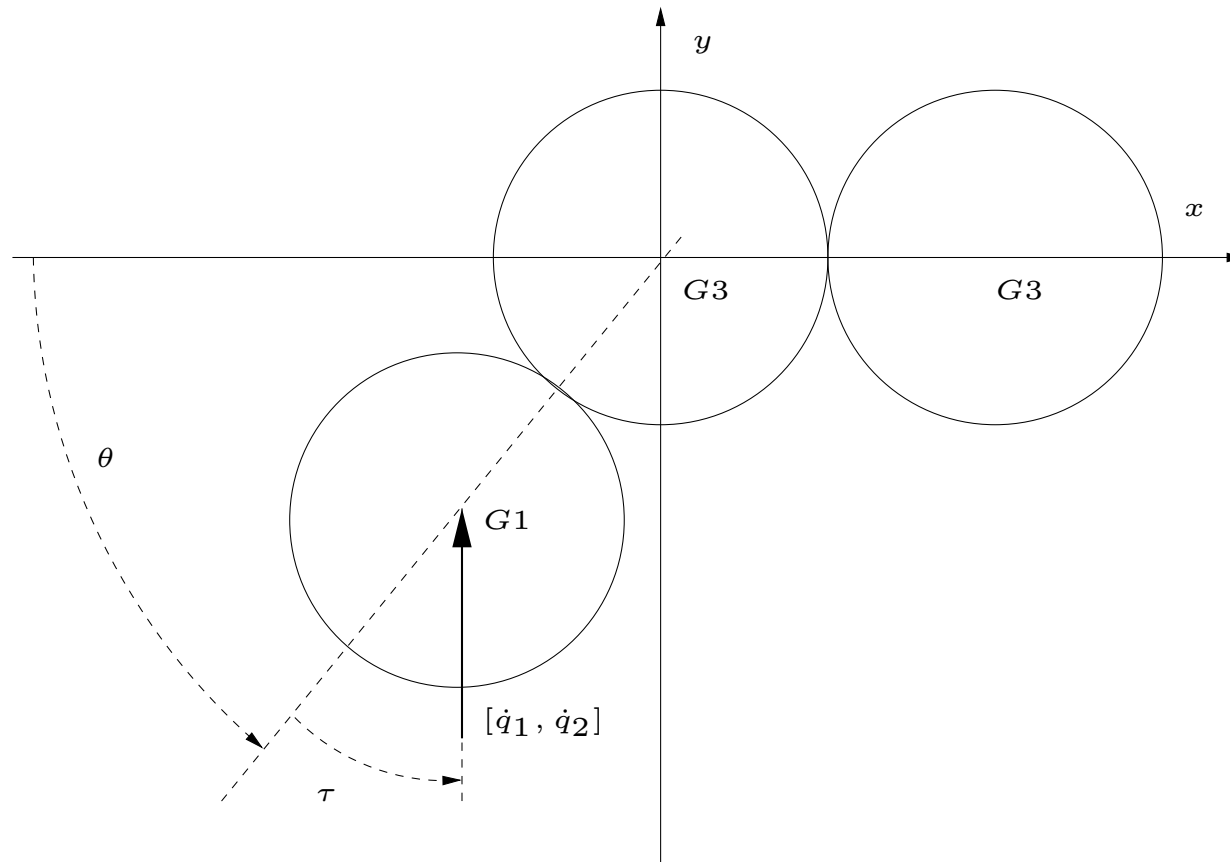
Major results:

- ✱ The ratio of impulses is finite and the subspace $E = \{\delta \geq 0, \dot{\delta} \geq 0\}$ is globally attractive.
- ✱ The impulse ratios are independent of the absolute value of stiffness and masses but only function of the ratio of stiffness and mass.
- ✱ In the linear case $\nu = 1$, the impulse ratios are completely determined by the natural modes of the regularized dynamical system and the pre-impact velocities.

Conclusion:

- ✱ The definition of the impulse ratio is relevant in the rigid limit.
- ✱ The impulse ratio provide us with some informations on the dynamical process
- ✱ The independence of the absolute value of stiffness avoids stiff numerical problem.

- The pool: A bidimensional chain of balls



3-balls in the plane

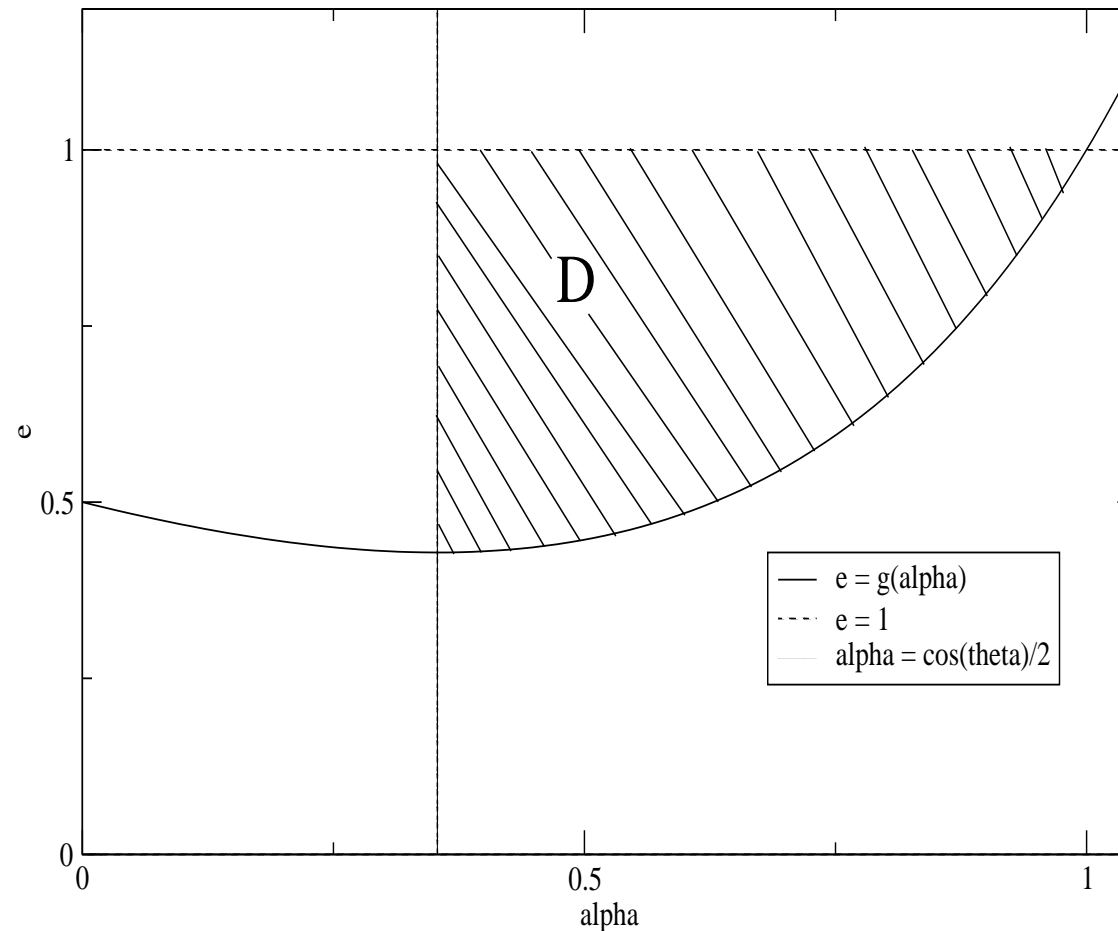
The pool: A bidimensional chain of balls

Solution for $\tau = 0$ and $v_1 = 1$

$$\left\{ \begin{array}{l}
 p_1 = \frac{1 + \sqrt{\Delta}}{-2\alpha_{12} \cos \theta + 2\alpha_{12}^2 + 2} \\
 p_2 = \frac{(1 + \sqrt{\Delta})\alpha_{12}}{-2\alpha_{12} \cos \theta + 2\alpha_{12}^2 + 2} \\
 \dot{x}_1^+ = \cos \theta \left(1 - \frac{1 + \sqrt{\Delta}}{-2\alpha_{12} \cos \theta + 2\alpha_{12}^2 + 2} \right) \\
 \dot{y}_1^+ = \sin \theta \left(1 - \frac{1 + \sqrt{\Delta}}{-2\alpha_{12} \cos \theta + 2\alpha_{12}^2 + 2} \right) \\
 \dot{x}_2^+ = (\cos \theta - \alpha_{12}) \left(\frac{1 + \sqrt{\Delta}}{-2\alpha_{12} \cos \theta + 2\alpha_{12}^2 + 2} \right) \\
 \dot{y}_2^+ = \sin \theta \left(\frac{1 + \sqrt{\Delta}}{-2\alpha_{12} \cos \theta + 2\alpha_{12}^2 + 2} \right) \\
 \dot{x}_3^+ = \alpha_{12} \left(\frac{1 + \sqrt{\Delta}}{-2\alpha_{12} \cos \theta + 2\alpha_{12}^2 + 2} \right) \\
 \dot{y}_3^+ = 0 \\
 \text{with} \\
 \Delta = 2(e - 1)(\alpha_{12}^2 - \alpha_{12} \cos \theta + 1) + 1
 \end{array} \right. \quad (13)$$

- Valid domain for e and α

- ✱ Bounds imposed by $\Delta > 0$ and unilateral constraints



Valid domain of e and α for $\theta = \frac{\pi}{4}$.

Influence of the configuration and the initial conditions

✱ Influence of the angle τ .

- The angle τ determines the occurrence of an impact (single or multiple).

$$-v_1^- \in \mathcal{T}_1 = \{p \mid -\nabla h_1(q) \cdot p \geq 0\}. \quad (13)$$

- The model leads to consistent results whatever the angle τ .

✱ Influence of the angle θ for τ such that $-v_1^- \in \mathcal{T}_1$:

1. $\theta \in [0, \frac{\pi}{2})$, one multiple impact

- We obtain a unique solution. The set of solutions is parametrized by e and α .

2. $\theta \in [\frac{\pi}{2}, \frac{2\pi}{3})$, two single impacts

- We must choose a correct value, i.e. $\alpha_{12} = 0$ to have consistent results.

or

- We must choose two standards law for single impact.

3. $\theta = \frac{2\pi}{3}$, The Bernoulli's problem

- A impulse ratio $\alpha_{13} = p_3/p_1$ is added.
- The choice $\alpha_{12} = \alpha_{13}$ leads to consistent results in the symmetric case (same masses and stiffnesses.).

• Preliminary conclusion

✱ Validity of the model.

- In the case of one multiple impact, the set of solutions which respect the equations of motion, the unilateral constraints and the energetic balance is correctly parametrized by the parameter e and α_{12} ,
- In the case of two single impacts, standard single laws must be applied or correct values of the impulse ratios must be chosen.

✱ Open problems

- How to separate multiple or single impact cases ?
- How to write a single model for all impacting cases ?

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- 3 – Lagrangian systems
 - 3.1 – General model
 - 3.2 – Angles between constraints
 - 3.3 – Illustration on an open chain of balls
- 4 – Conclusion and perspectives

General model

- ✱ equations of motion at the instant of impact.

$$M(\dot{q}^+ - \dot{q}^-) = P, \quad q \in \mathbb{R}^n \quad (14)$$

- ✱ unilateral constraints

$$h_i(q) \geq 0, \quad i = 1, \dots, m \quad (15)$$

- ✱ kinematic relations :

$$v = H\dot{q} \quad (16)$$

$$P = H^T p \quad (17)$$

$$H = \nabla h(q(t_k)) \quad (18)$$

- ✱ unilateral constraints on the relative velocity and the percussions

$$\text{if } h_i(q) = 0, \quad v_i^+ \geq 0, p_i \geq 0 \quad (19)$$

$$\text{if } h_i(q) > 0, \quad p_i = 0 \quad (20)$$

Angles between constraints in the kinetic metric

- ✱ Gradient to a constraint in the kinetic metric :

$$n_i = \frac{M^{-1}(q)\nabla_q h_i(q)}{\sqrt{\nabla_q h_i(q)^T M(q)^{-1} \nabla_q h_i(q)}} \quad (21)$$

- ✱ Angle between two constraints in the kinetic metric:

$$\cos \theta_{ij} = \langle n_i, n_j \rangle_M = \frac{\nabla_q^T h_i(q) M^{-1}(q) \nabla_q h_j(q)}{\sqrt{\nabla_q h_i(q)^T M(q)^{-1} \nabla_q h_i(q)} \sqrt{\nabla_q h_j(q)^T M(q)^{-1} \nabla_q h_j(q)}} \quad (22)$$

It is noteworthy that $\nabla_q^T h_i(q) M^{-1}(q) \nabla_q h_j(q)$ is a coefficient of the Delassus matrix $H^T M^{-1} H$

- ✱ Coupling between two adjacent constraints:

$$\cos \theta_{ij} = \langle n_i, n_j \rangle_M < 0 \implies \text{coupling } \alpha_{ij} = p_i/p_j \quad (23)$$

$$\cos \theta_{ij} = \langle n_i, n_j \rangle_M \geq 0 \implies \text{no coupling due to unilaterality} \quad (24)$$

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• Illustration on an open chain of balls

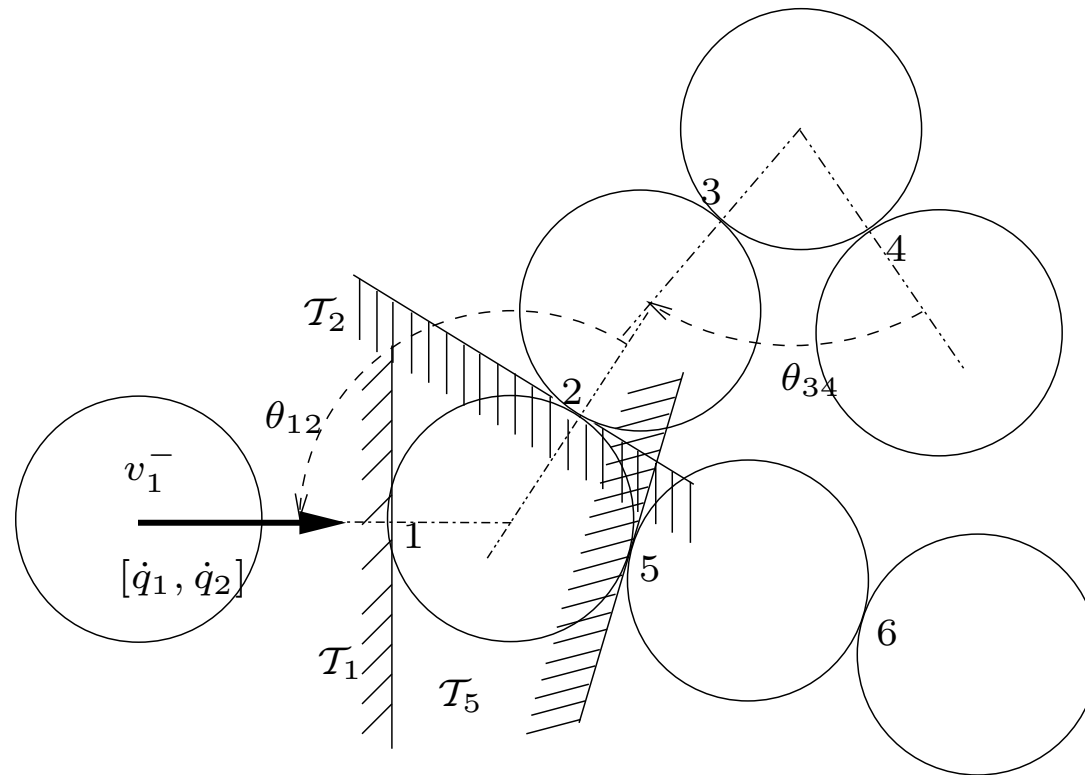


Illustration on an open chain of balls

- ✱ First criteria of impact (single or multiple), the velocity v_1 must belong to the opposite of the tangent cone $\mathcal{T}_1(q)$

$$-v_1^- \in \mathcal{T}_1 = \{p \mid -\nabla h_1(q) \cdot p \geq 0\}. \quad (26)$$

- ✱ To have a multiple impact of degree k , (involving the first k -balls) the $k - 1$ first kinetic angle $\theta_{i,i+1}$ must be greater than $\frac{\pi}{2}$:

$$\theta_{i,i+1} > \frac{\pi}{2}, \text{ i.e. } \langle n_i, n_{i+1} \rangle_M < 0, i = 1 \dots k - 1 \quad (25)$$

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Conclusion and perspectives

- ✱ Model with a global energetic coefficient and impulse ratios is a good candidate for propagation of an impact
- ✱ Properties of the impulse ratio
 - Well defined in the rigid limit of a regularized model
 - Easy to evaluate from the numerical point of view.
 - Parametrization of the set of solutions after impact.
- ✱ A general impact law for multiple and single impacts is a more challenging task:
 - How to take into account the kinetic angle between constraints into the formulation ?
 - How to write a single law valid in all cases ?
 - How to deal with such law on the numerical point of view ?

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